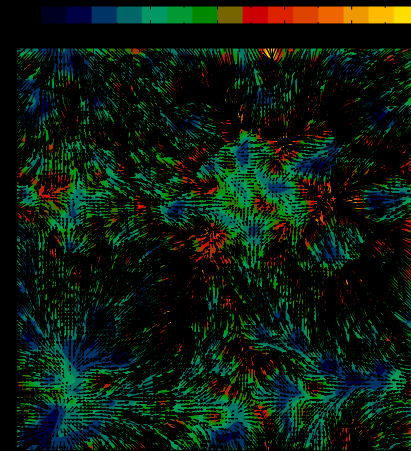
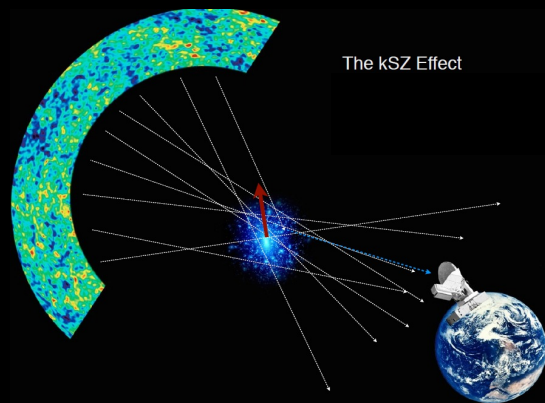
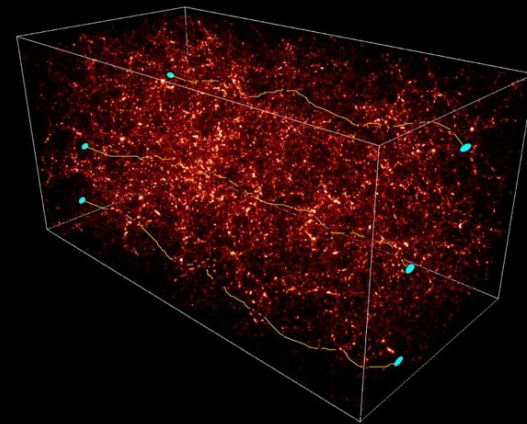
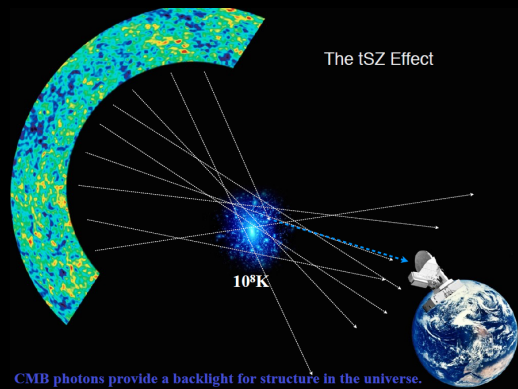


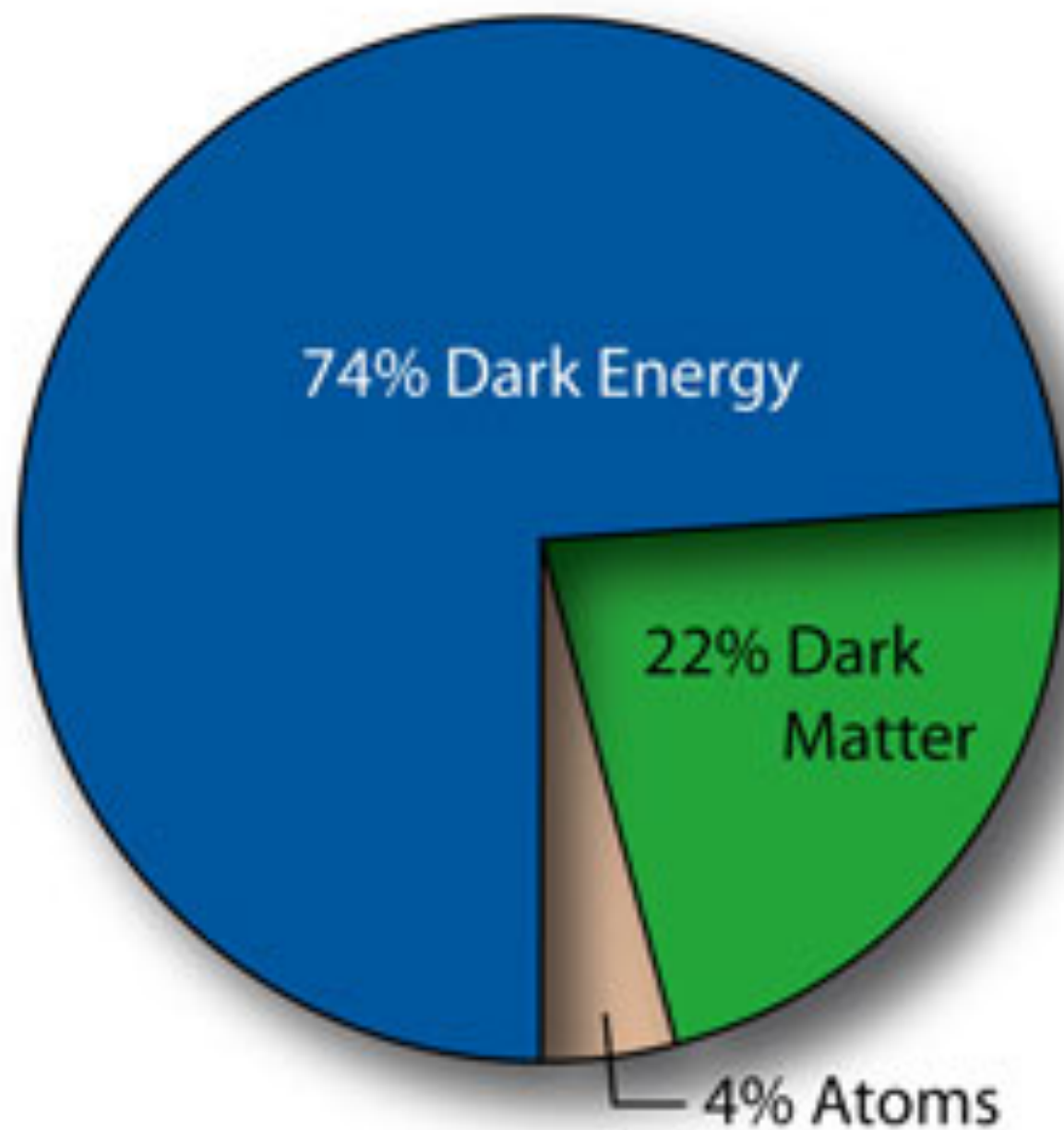
# Detection of the Missing Baryons

## Yin-Zhe Ma (马寅哲)



Postdoc at Jodrell Bank, Manchester (2014-2015, C.Dickinson R.Battye)

→ University of KwaZulu-Natal, Durban, South Africa (Sep 2015 --)

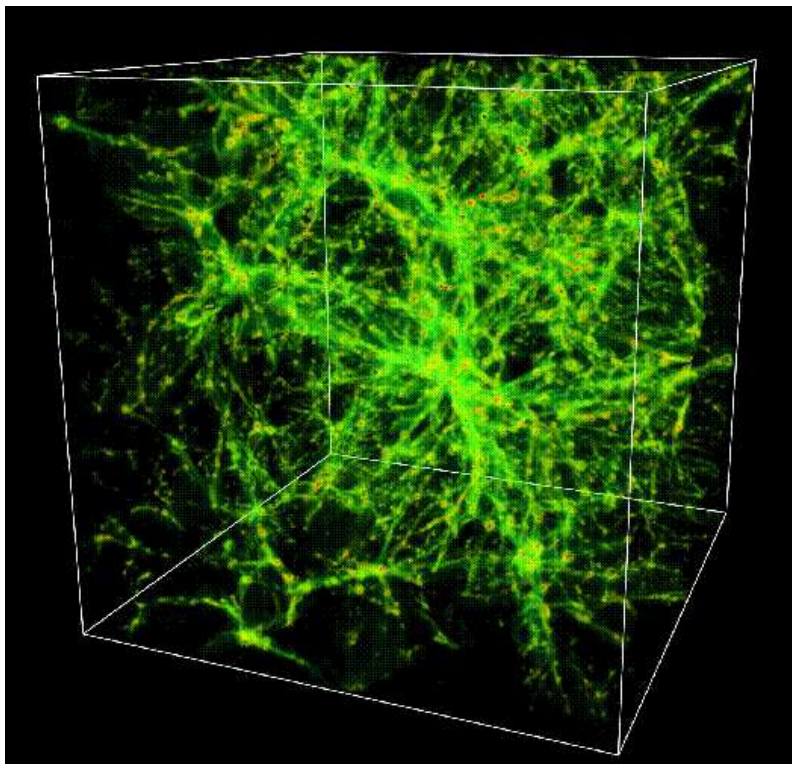


# Cosmic baryon inventory:

Category	Parameter	Components <sup>a</sup>
3.3.....	Main-sequence stars: spheroids and bulges	$0.0015 \pm 0.0004$
3.4.....	Main-sequence stars: disks and irregulars	$0.00055 \pm 0.00014$
3.5.....	White dwarfs	$0.00036 \pm 0.00008$
3.6.....	Neutron stars	$0.00005 \pm 0.00002$
3.7.....	Black holes	$0.00007 \pm 0.00002$
3.8.....	Substellar objects	$0.00014 \pm 0.00007$
3.9.....	H I + He I	$0.00062 \pm 0.00010$
3.10.....	Molecular gas	$0.00016 \pm 0.00006$
3.11.....	Planets	$10^{-6}$
3.12.....	Condensed matter	$10^{-5.6 \pm 0.3}$
3.13.....	Sequestered in massive black holes	$10^{-5.4}(1 + \epsilon_n)$

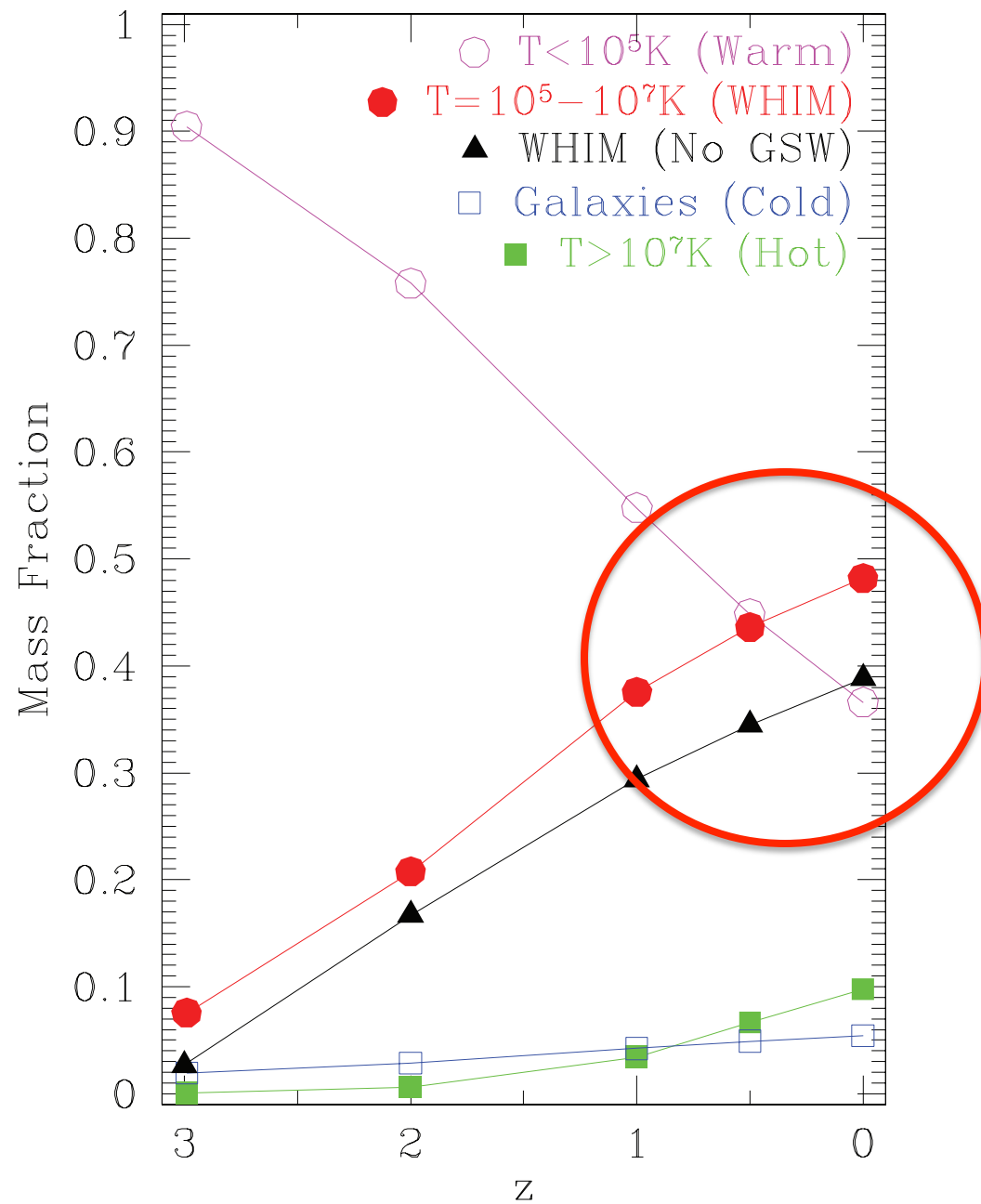
3.3+...+3.13:  $\Omega_{b,g} = 0.0035$  =8% total baryon density

90% of baryons are in either intergalactic or intercluster medium



Cen and Ostriker 2006

X-ray:  $\sim n_e(r)^2$





thermal Sunyaev-Zeldovich effect

X

Weak Lensing

and compare with halo model prediction and hydrodynamic simulation

kinetic Sunyaev-Zeldovich effect

X

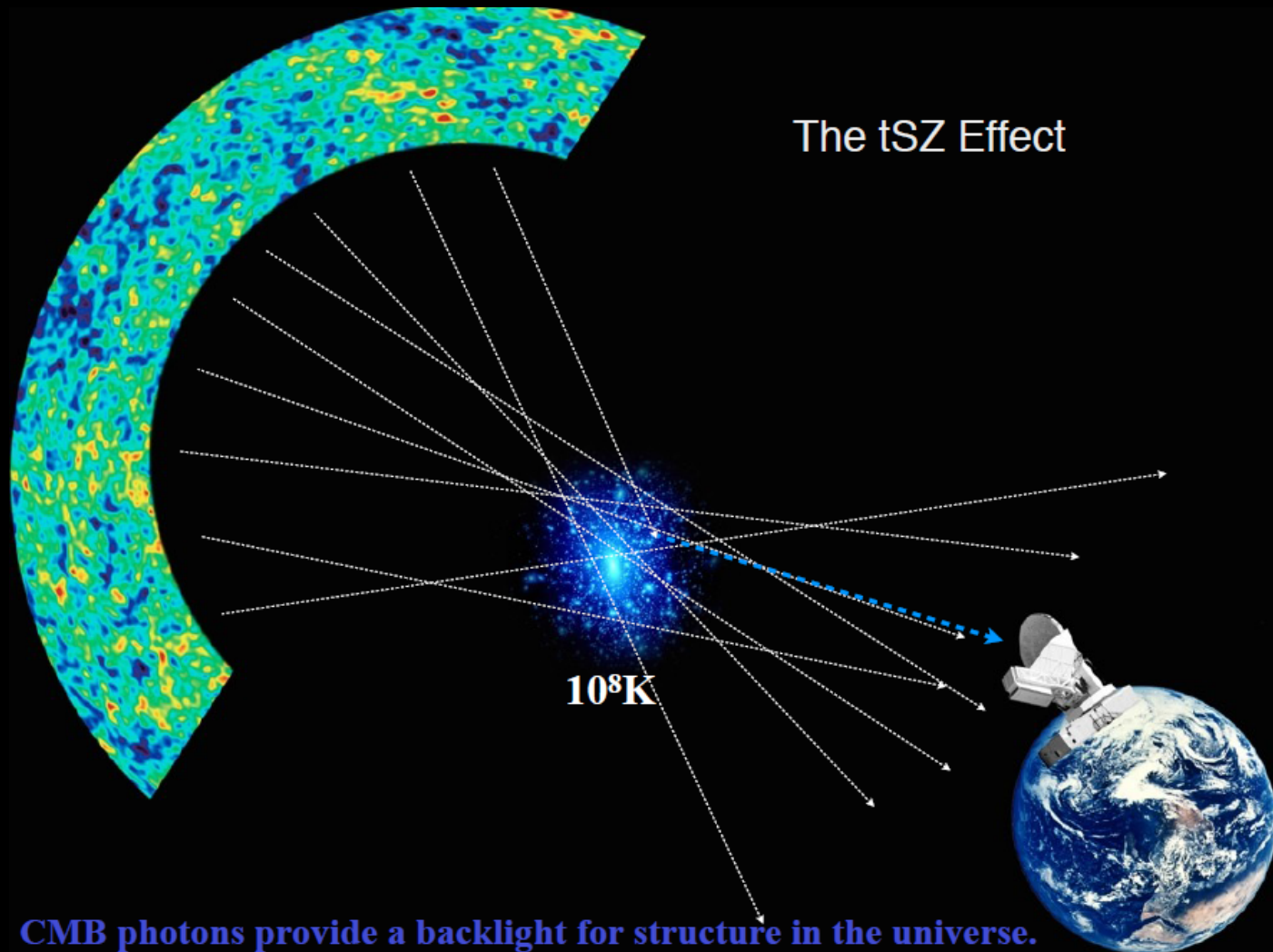
Peculiar velocity field

Thermal SZ maps

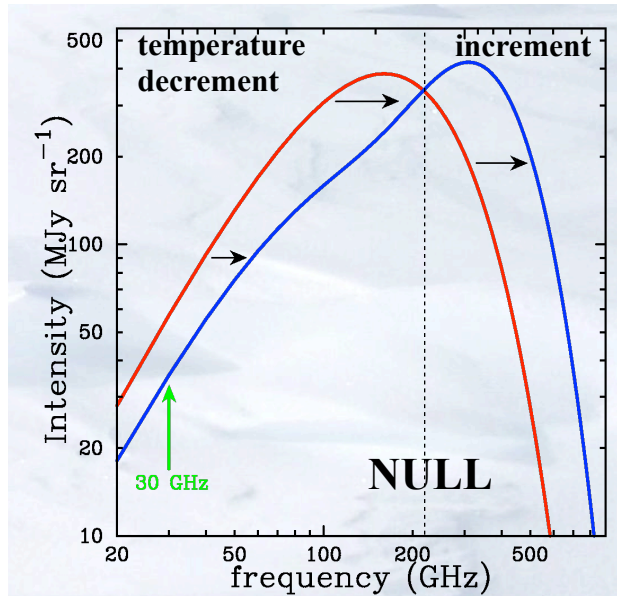
X

Luminous red galaxies

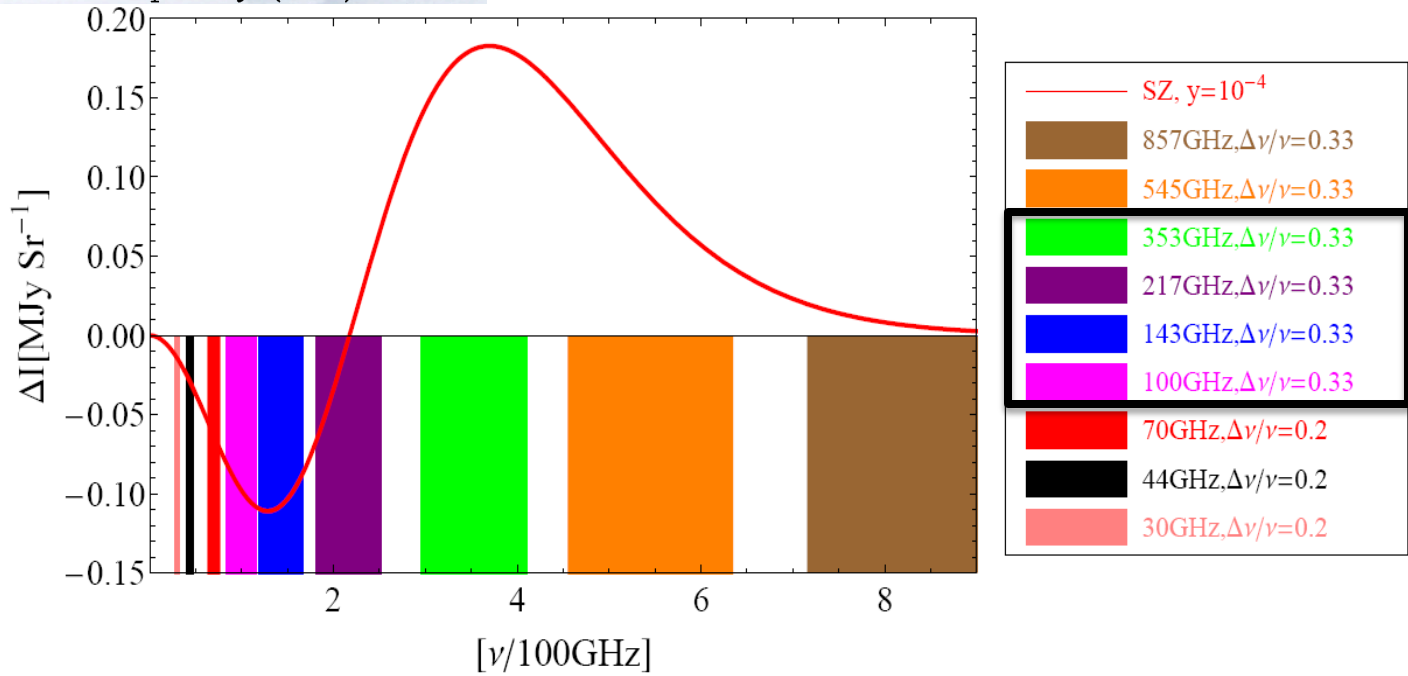
# The thermal Sunyaev-Zeldovich effect



# Thermal Sunyaev-Zeldovich effect (tSZ):



$$y = \frac{k_B \sigma_T}{m_e c^2} \int_0^l T_e(l) n_e(l) dl$$



SZ map from linear combination of Planck frequency bands:  
 $\nu_i = 100, 143, 217, 353$  GHz.

$$T_{SZ}/T_0 \equiv y S_{SZ}(\nu_i) = \sum b_i T(\nu_i)$$

- $\sum b_i S_{SZ}(\nu_i) = 1$        $S_{SZ}(x) = x \coth(x/2) - 4$  ( $x = h\nu/kT$ )
- $\sum b_i S_{CMB}(\nu_i) = 0$        $S_{CMB}(x) = 1$
- $\sum b_i S_{\text{dust}}(\nu_i) = 0$        $S_{\text{dust}}(\nu_i) = \nu^\beta g(x)$

# Planck Full-Sky Maps – 4 Frequencies

30 GHz

44 GHz

70 GHz

100 GHz

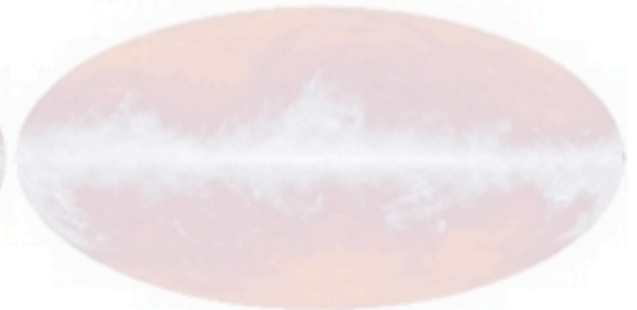
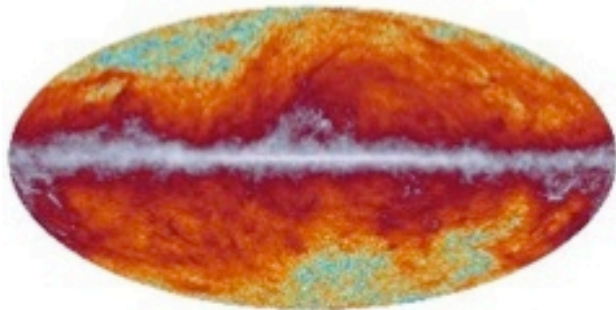
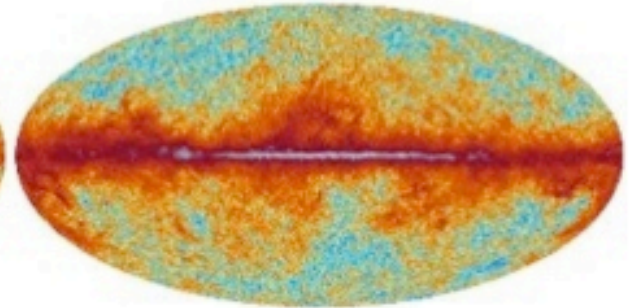
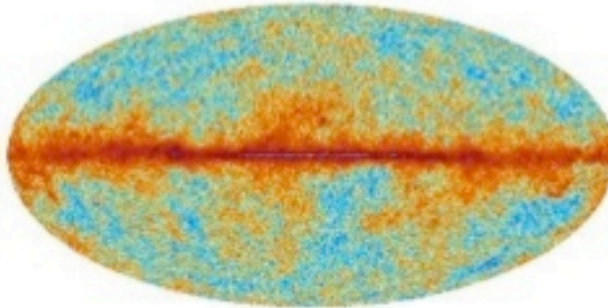
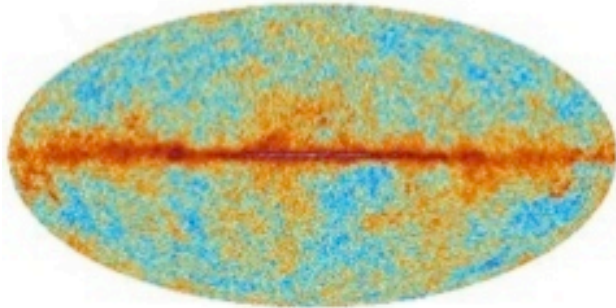
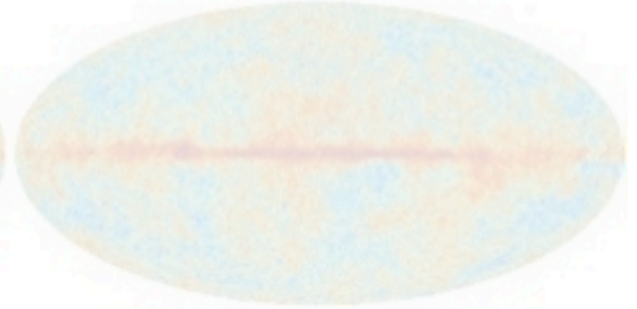
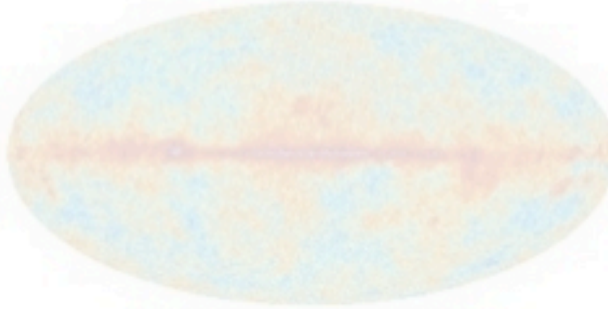
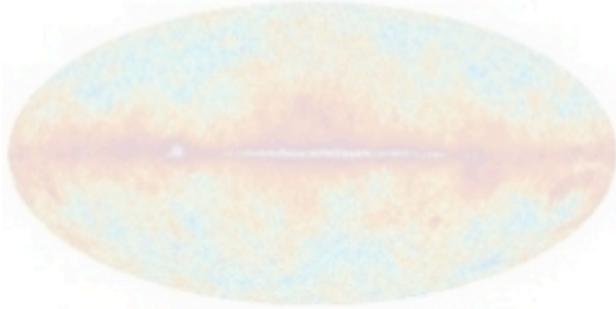
143 GHz

217 GHz

353 GHz

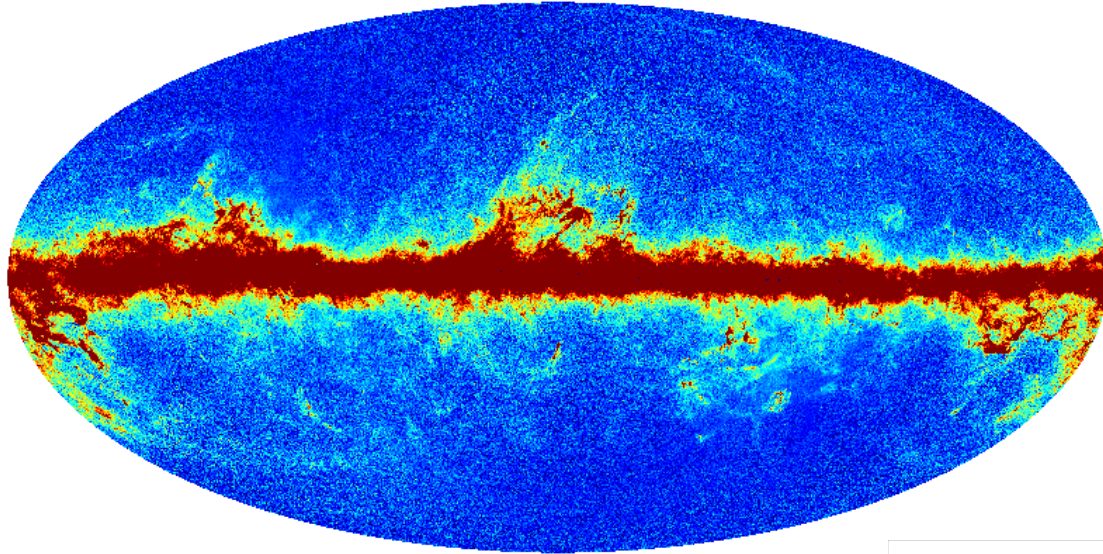
545 GHz

857 GHz





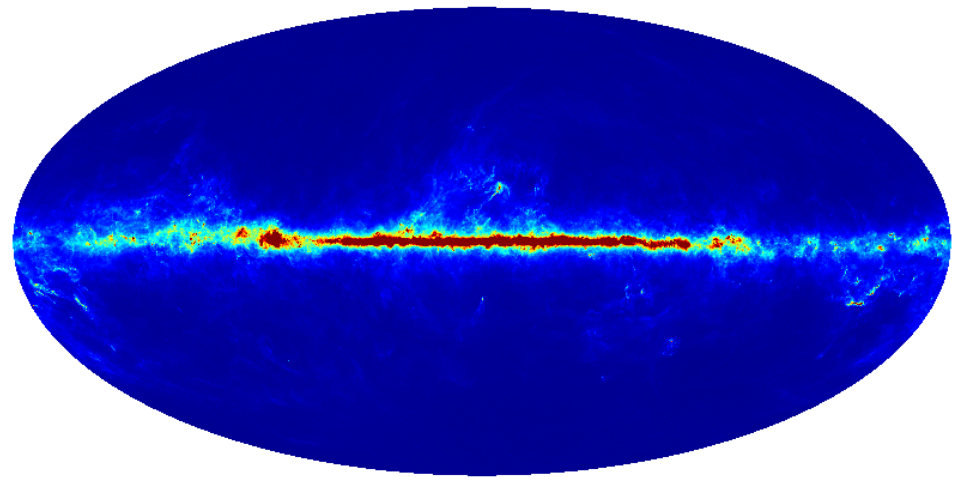
# Planck SZ y map, version E



y=0   $10^{-4}$

Reject  $\beta_{\text{dust}} = 2.0$ ,  $r_{2.0}(100 \text{ GHz}) = 0$

# Planck SZ no-y map, version E



y=0   $10^{-4}$

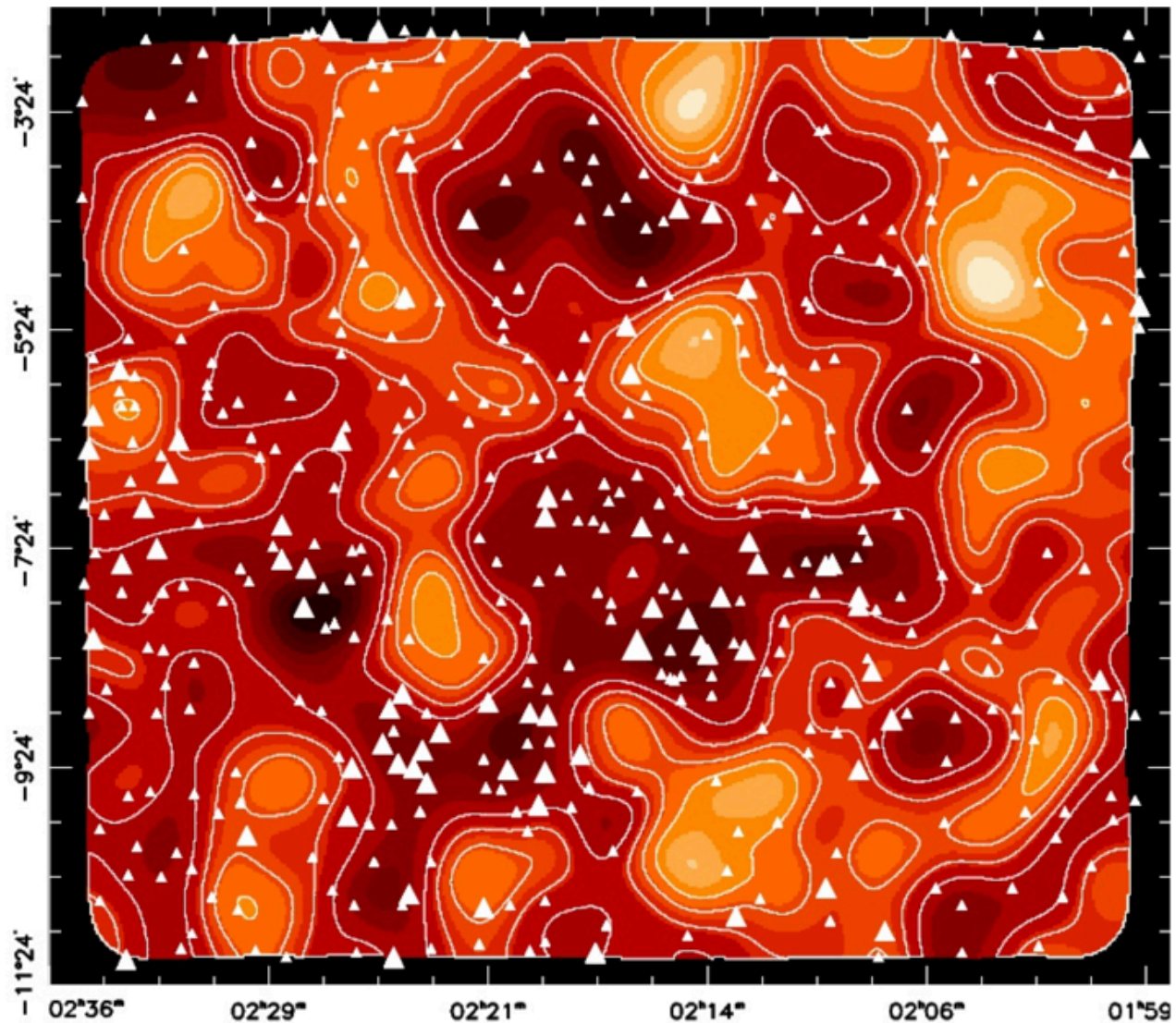
Reject  $S_{\text{SZ}}(\nu)$ , retain  $\beta_{\text{dust}} = 1.8$

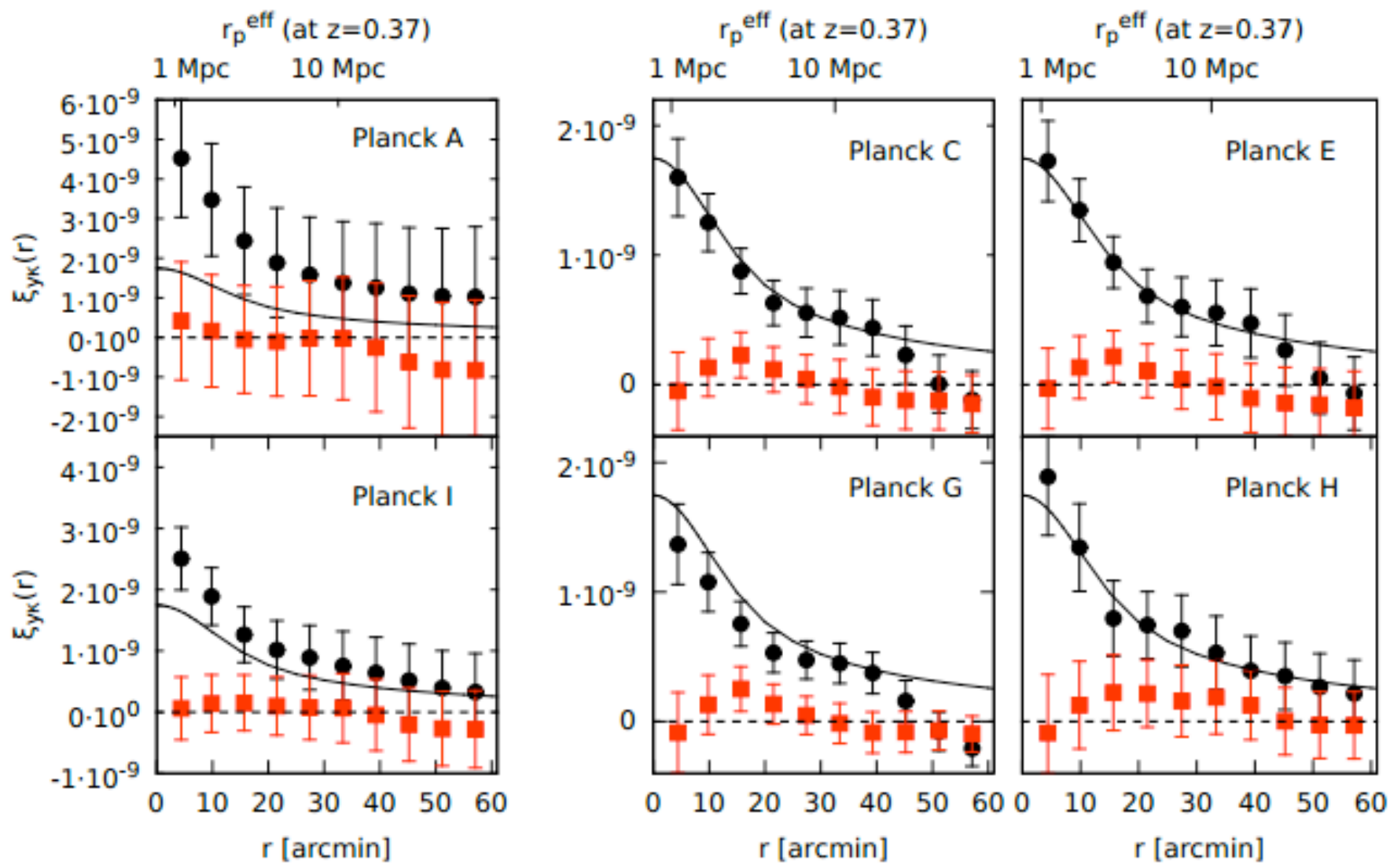


# CFHT mass map:

154 deg<sup>2</sup> in 4  
patches

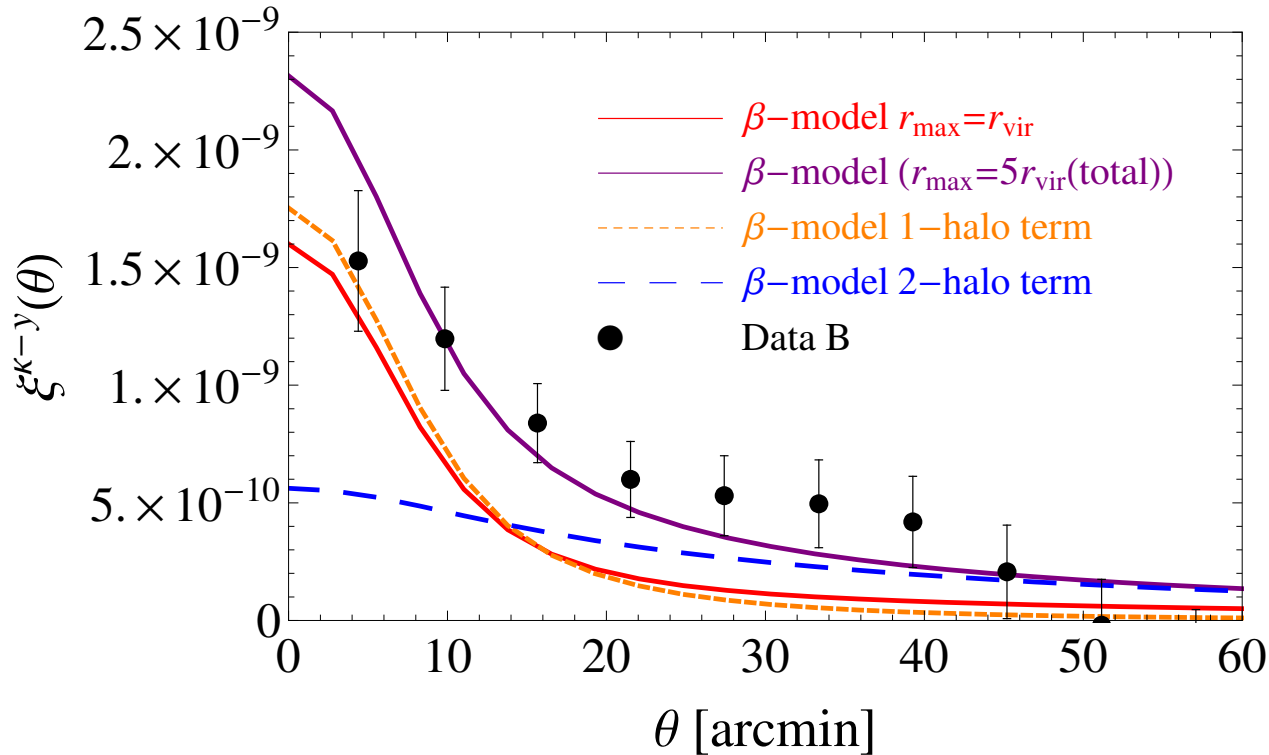
Van Waerbeke et  
al., 2014, MNRAS





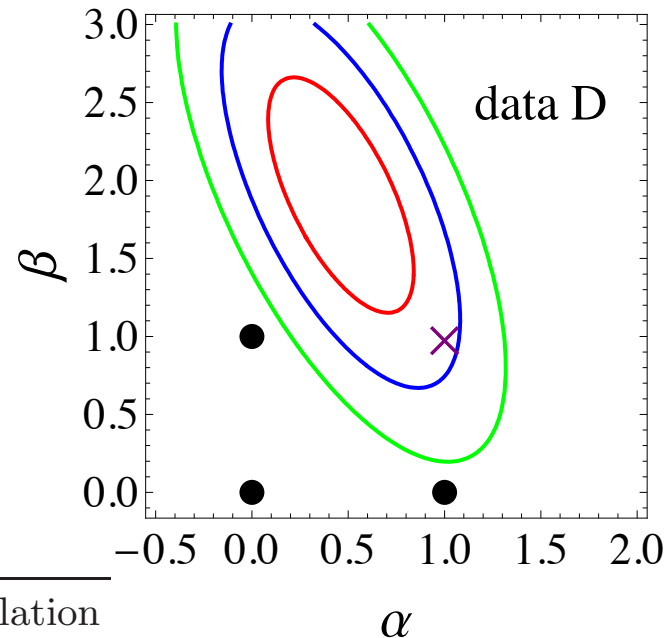
van Waerbake, Hinshaw, Murray: 2014 Phys. Rev. D

# Halo model:



Ma et al. fits a halo model to the observed correlation function. A  $\beta$  model fits well, but in this context the data requires a 2-halo term to fit the large angular scale separation.

$$\chi^2(\alpha, \beta) = \sum_{ij} [\xi^d(\theta_i) - \alpha\xi^{1h}(\theta_i) - \beta\xi^{2h}(\theta_i)] \\ \times C_{ij}^{-1} [\xi^d(\theta_j) - \alpha\xi^{1h}(\theta_j) - \beta\xi^{2h}(\theta_j)]$$



Data set	2-halo only	1-halo only	No correlation
B	$7.6 \times 10^{-5}$	$2.4 \times 10^{-4}$	$3.8 \times 10^{-11}$
C	$2.2 \times 10^{-5}$	$1.0 \times 10^{-4}$	$1.1 \times 10^{-11}$
D	$7.1 \times 10^{-5}$	$6.0 \times 10^{-4}$	$6.6 \times 10^{-11}$
E	$2.6 \times 10^{-5}$	$2.8 \times 10^{-4}$	$1.3 \times 10^{-11}$
F	$1.7 \times 10^{-3}$	$5.4 \times 10^{-3}$	$1.5 \times 10^{-8}$
G	$4.6 \times 10^{-3}$	$1.0 \times 10^{-2}$	$9.6 \times 10^{-8}$
H	$6.7 \times 10^{-4}$	$7.3 \times 10^{-5}$	$1.1 \times 10^{-9}$

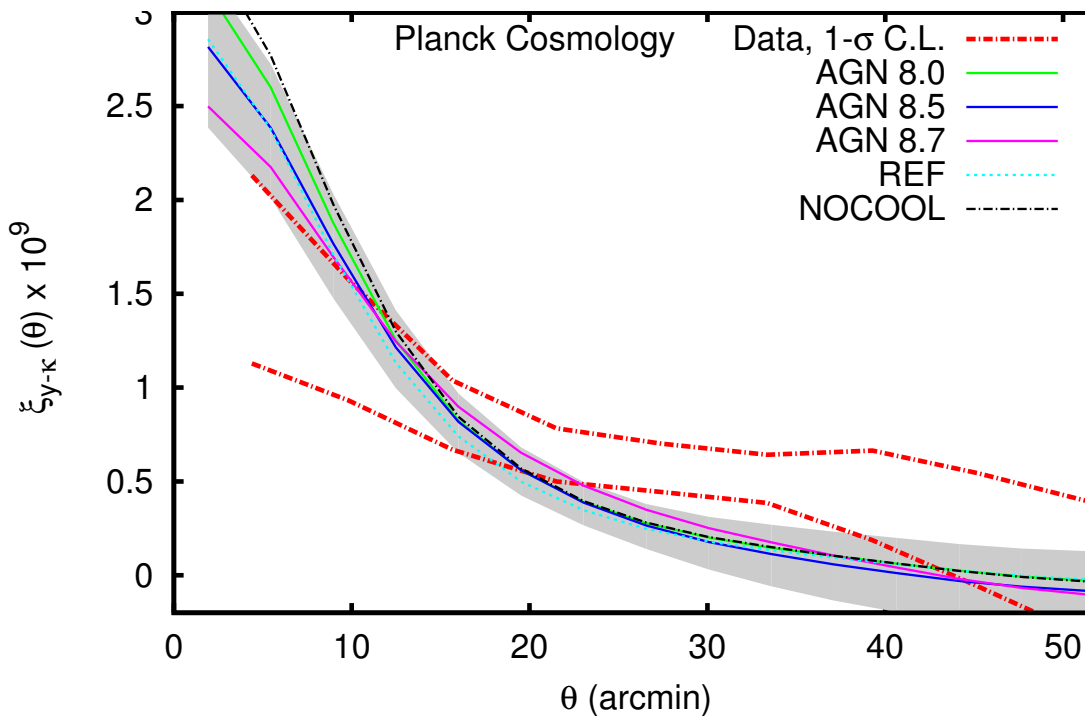
TABLE I: For each  $y$ -map B–H, the probability that the fit in Eq. (6) allows:  $\alpha = 0, \beta = 1$  (no 1-halo term, column 2);  $\alpha = 1, \beta = 0$  (no 2-halo term, column 3); and  $\alpha = \beta = 0$  (no cross-correlation, column 4). We assume  $P = \exp(-\Delta\chi^2/2)$ .

	$10^{12} M_{\odot} - 10^{14} M_{\odot}$	$10^{14} M_{\odot} - 10^{16} M_{\odot}$
$(0.01-1) r_{\text{vir}}$	26%	28%
$(1-100) r_{\text{vir}}$	14%	32%

By applying the virial theorem with  $z = 0.37$ , for the mass range  $10^{12} - 10^{16} M_{\text{sun}}$ , we get  $T_e = 10^5 - 10^8 \text{ K}$ .

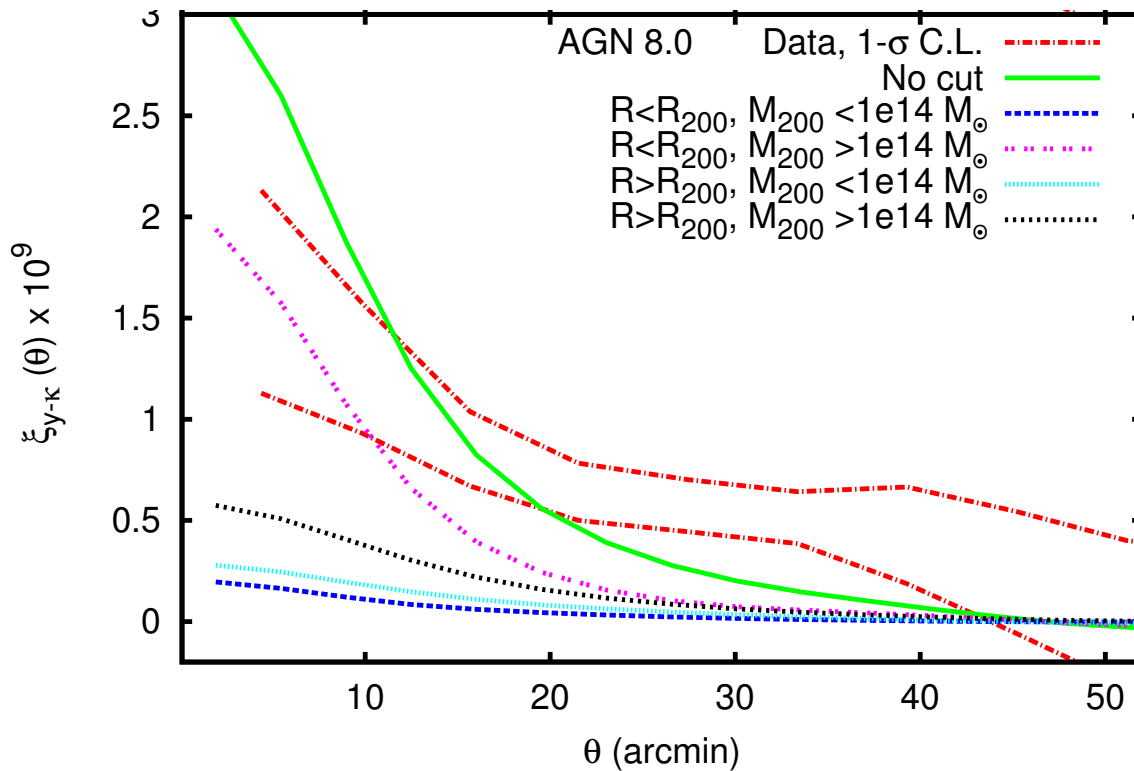
# Simulation vs data

Simulation	UV/X-ray background	Cooling	Star formation	SN feedback	AGN feedback	$\Delta T_{\text{heat}}$
NOCOOL	Yes	No	No	No	No	...
REF	Yes	Yes	Yes	Yes	No	...
AGN 8.0	Yes	Yes	Yes	Yes	Yes	$10^{8.0}$ K
AGN 8.5	Yes	Yes	Yes	Yes	Yes	$10^{8.5}$ K
AGN 8.7	Yes	Yes	Yes	Yes	Yes	$10^{8.7}$ K



NOCOOL is a standard non-radiative ('adiabatic') model. REF is the OWLS reference model and includes sub-grid prescriptions for star formation, metal-dependent radiative cooling, stellar evolution, mass loss, chemical enrichment, and a kinetic supernova feedback prescription. The AGN models are built on the REF model and additionally include a prescription for black hole growth and feedback from active galactic nuclei. The three AGN models differ only in their choice of the key parameter of the AGN feedback model  $\Delta T$ , which is the temperature by which neighboring gas is raised due to feedback.





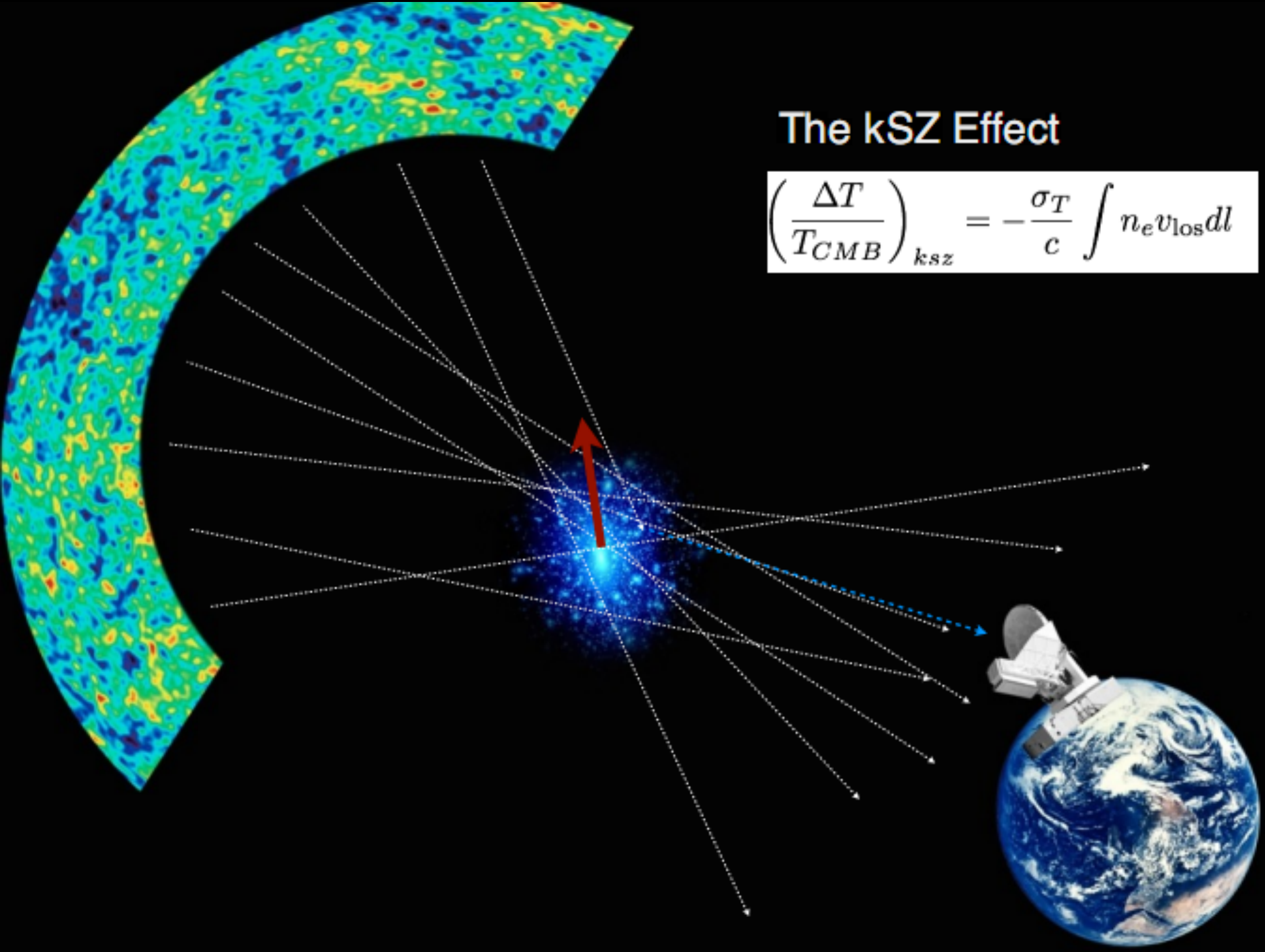
Bin	Simulation		Halo model
	Signal	Baryon	Signal
Low mass, inner radii	7%	6%	26%
Low mass, outer radii	12%	24%	14%
High mass, inner radii	56%	4%	28%
High mass, outer radii	25%	7%	32%

# Evidence of gas outside the virial radius

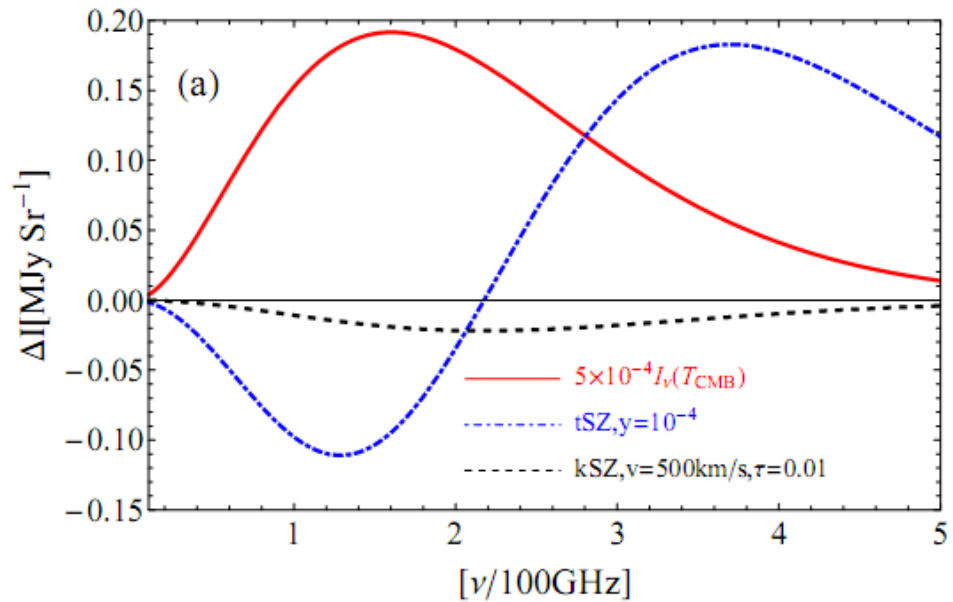
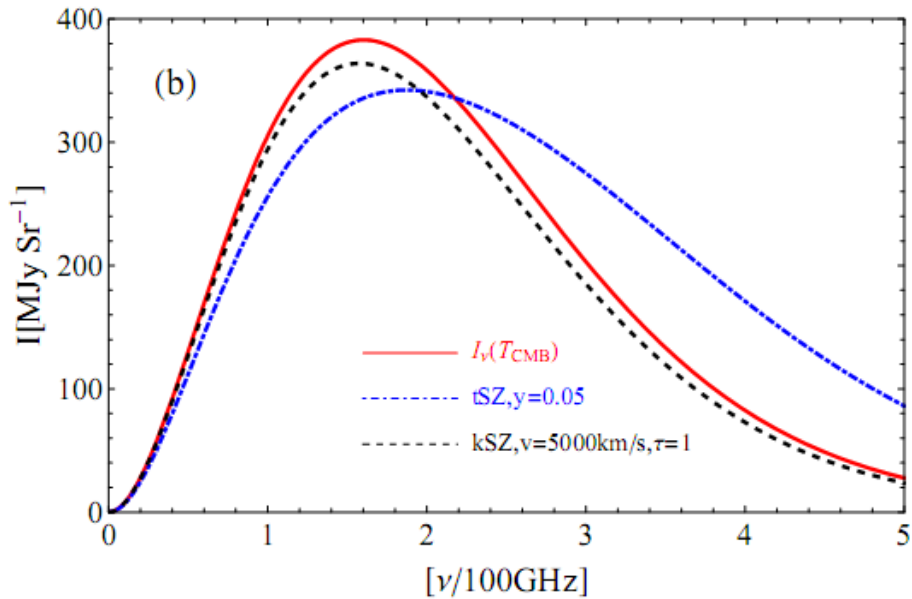
Planck intermediate results XXXVII, 1504.03339;  
C.Hernandez-Monteagudo, YZM, F-S Kitaura, W.Wang et al., 1504.04011

## The kSZ Effect

$$\left(\frac{\Delta T}{T_{CMB}}\right)_{kSZ} = -\frac{\sigma_T}{c} \int n_e v_{los} dl$$



kSZ: 
$$\frac{\delta T}{T_0}(\hat{\mathbf{n}}) = - \int dl \sigma_T n_e \left( \frac{\mathbf{v}}{c} \cdot \hat{\mathbf{n}} \right)$$



SZ map from linear combination of Planck frequency bands:  
 $\nu_i = 100, 143, 217, 353$  GHz.

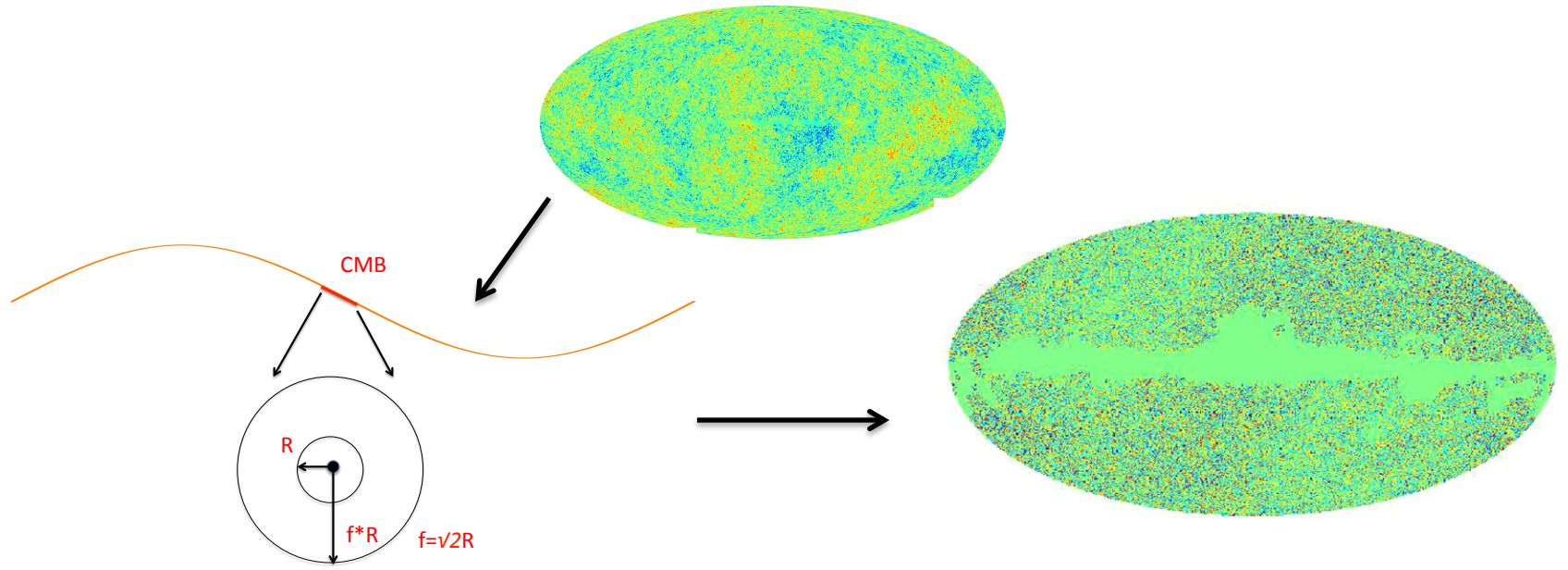
$$T_{SZ}/T_0 \equiv y S_{SZ}(\nu_i) = \sum b_i T(\nu_i)$$

$$1. \sum b_i S_{SZ}(\nu_i) = 1 \rightarrow 0! \quad S_{SZ}(x) = x \coth(x/2) - 4 \quad (x = h\nu/kT)$$

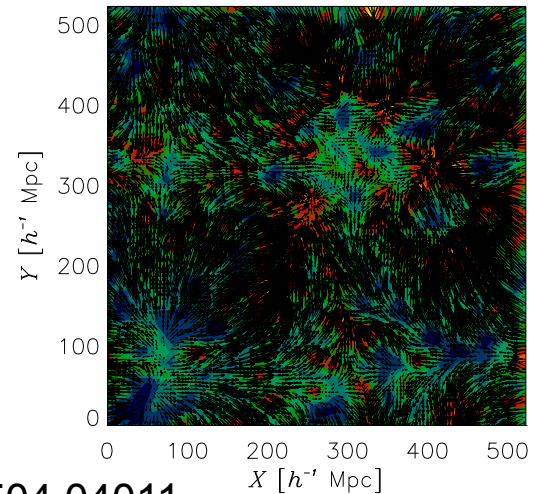
$$2. \sum b_i S_{CMB}(\nu_i) = 0 \rightarrow 1! \quad S_{CMB}(x) = 1$$

$$3. \sum b_i S_{\text{dust}}(\nu_i) = 0 \quad S_{\text{dust}}(\nu_i) = \nu^\beta g(x)$$

# Planck SMICA, SEVEM, NILC maps



**X**

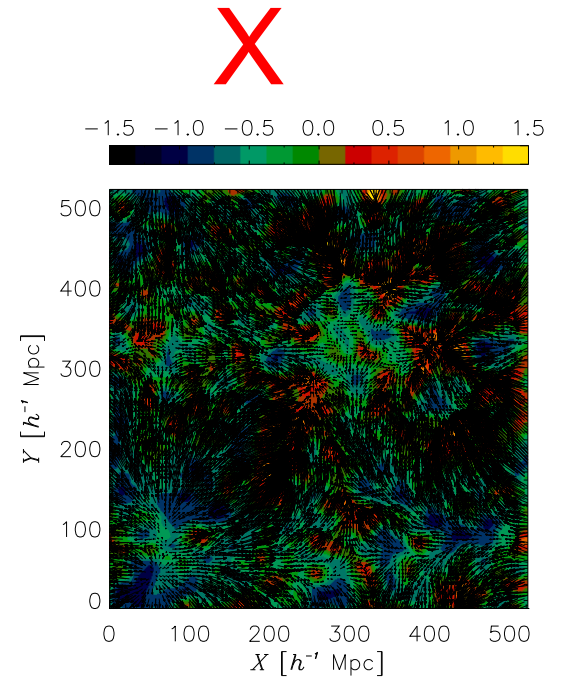
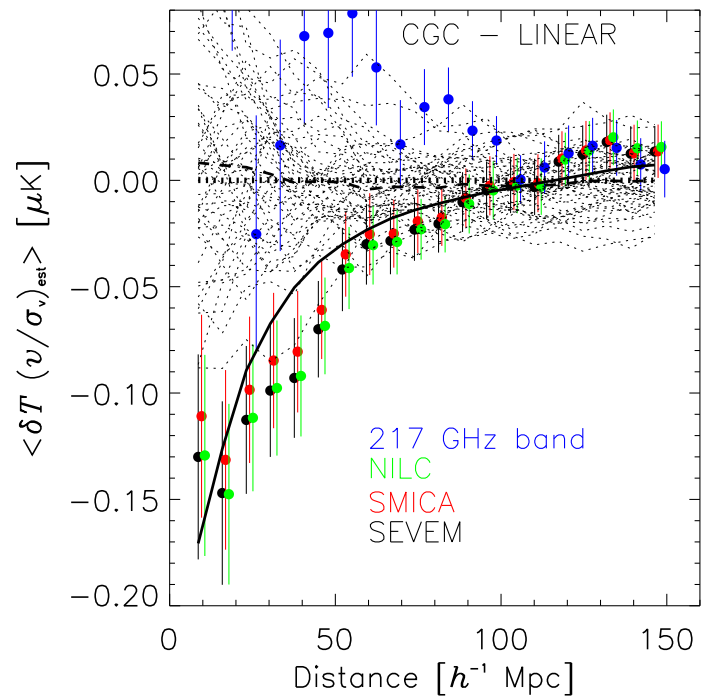
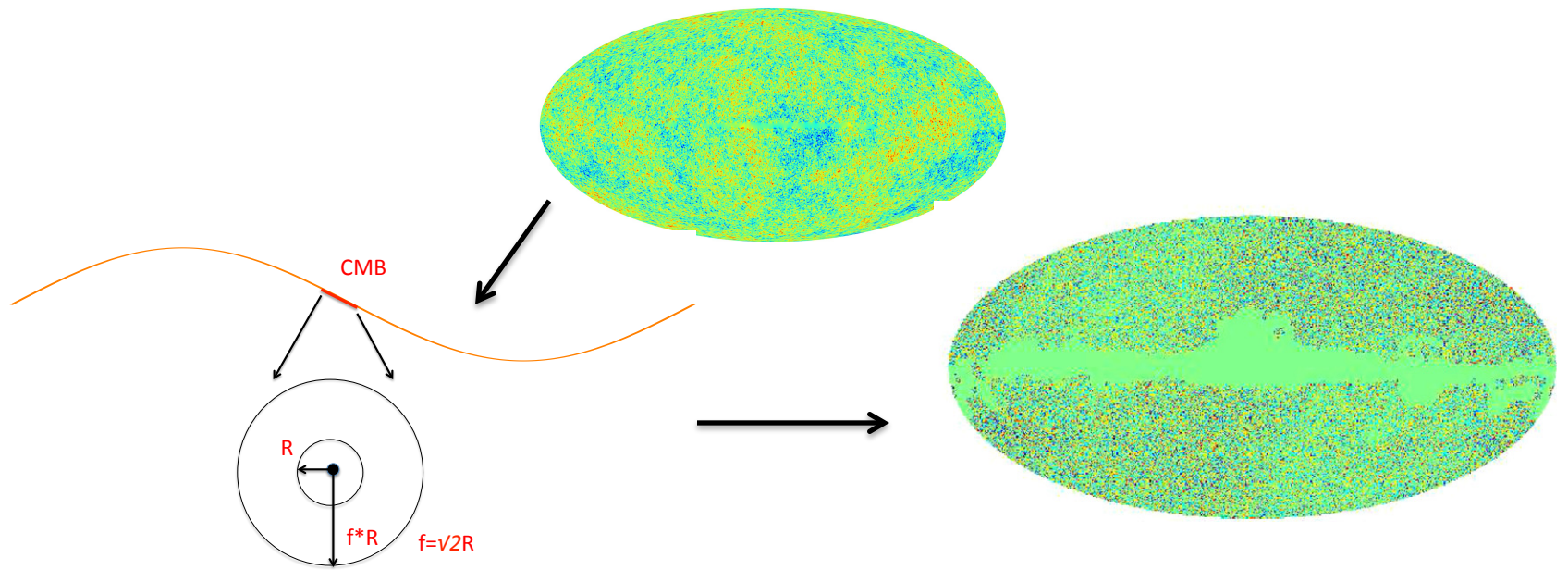


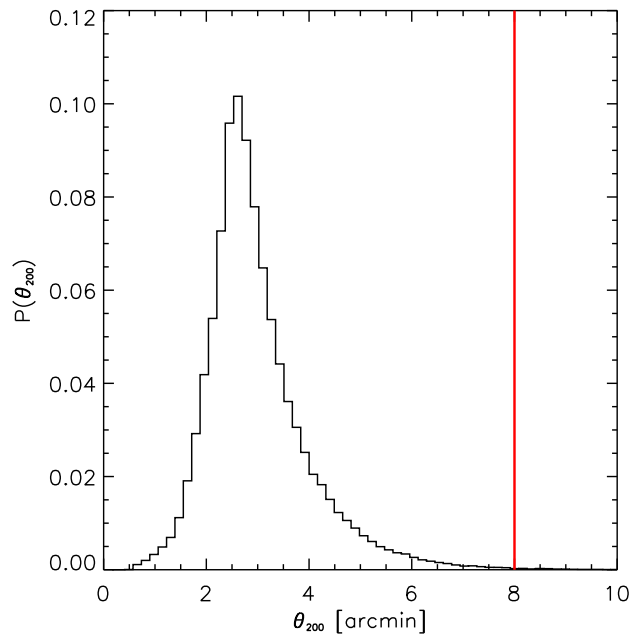
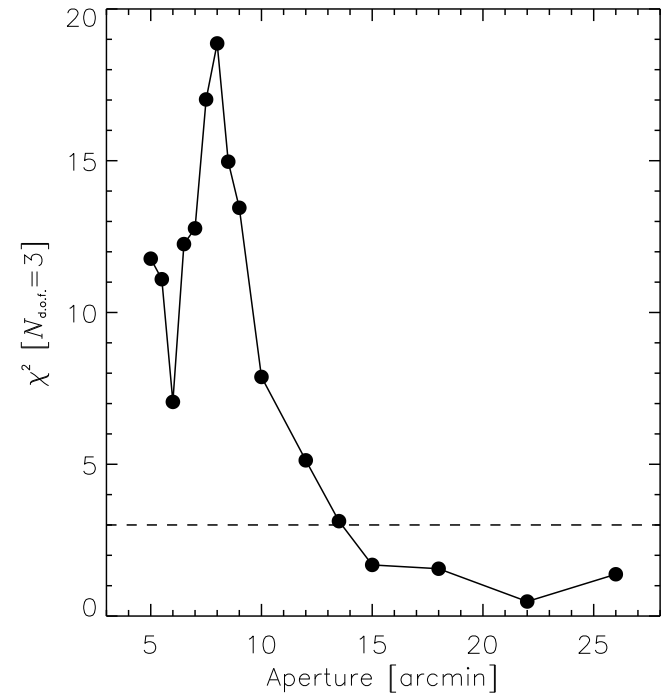
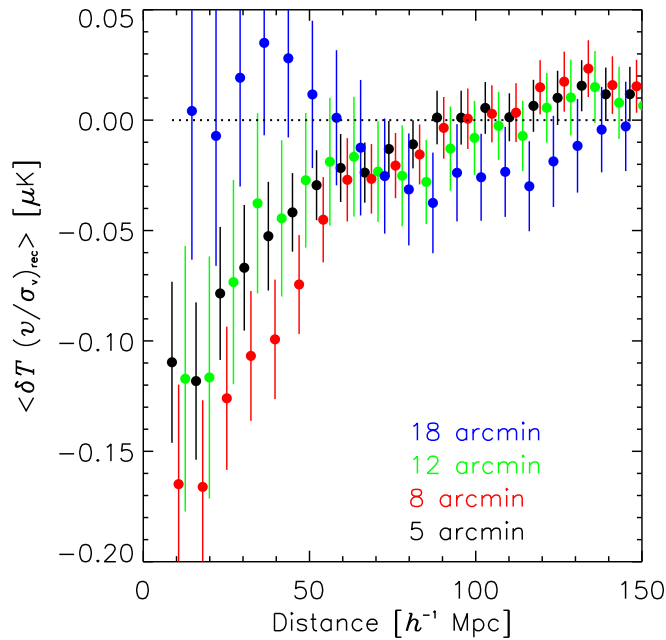
$$w^{T,v}(r) = \langle \delta T_i v_{\text{los}}^{\text{rec}}(\mathbf{x}_j) \rangle_{i,j}(r)$$





# Planck SMICA, SEVEM, NILC maps

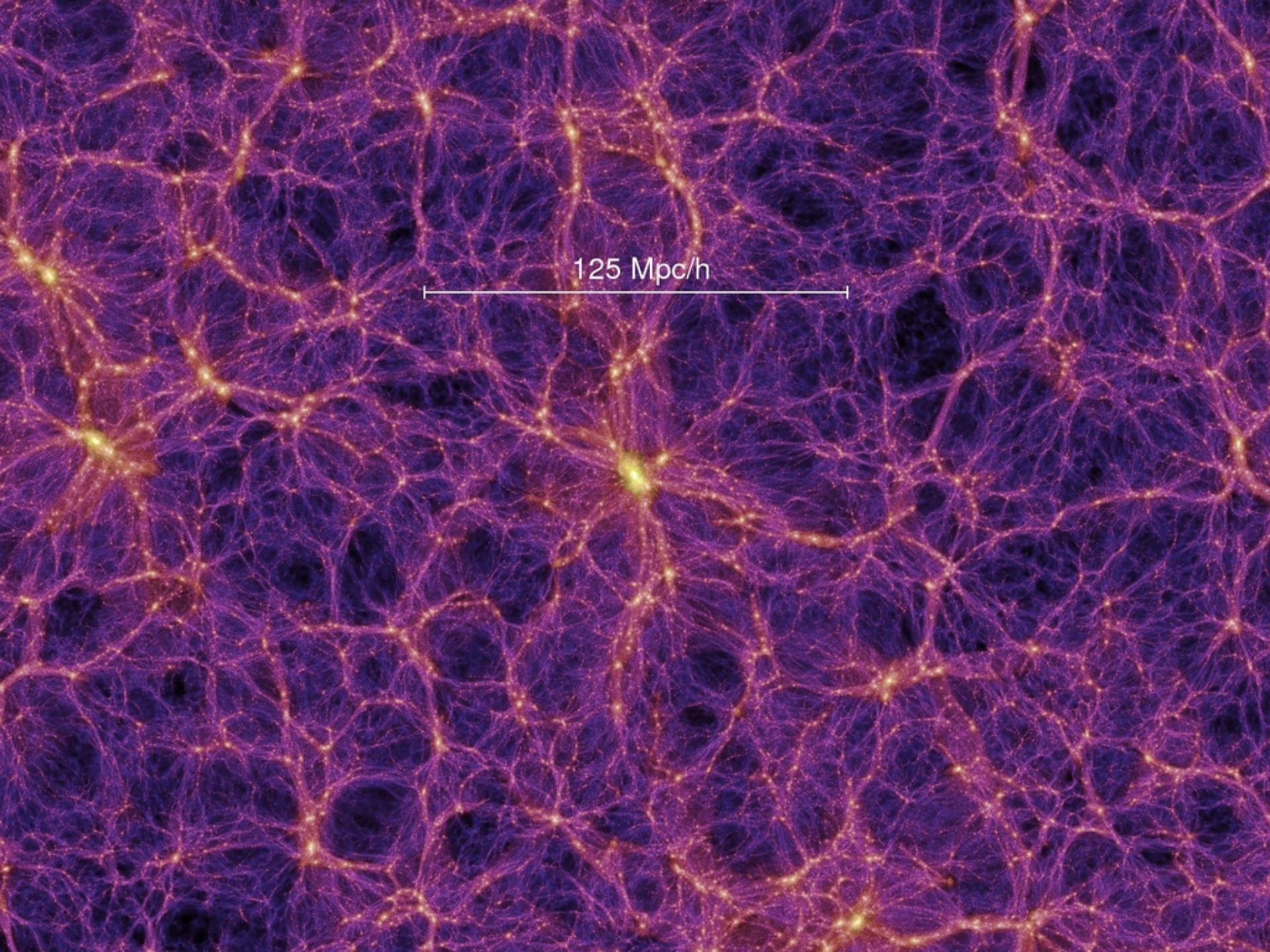




An aperture of 8 arcmin for CGC sources (placed at a median redshift of  $z = 0.12$ ) corresponds to a physical radius of roughly 1 Mpc, which is at least twice the typical virial radius of the CGs. Combining this with the behaviour of the signal for the 5, 8, and 12 arcmin apertures suggests that most of the detected kSZ signal is coming from outside the halos themselves.

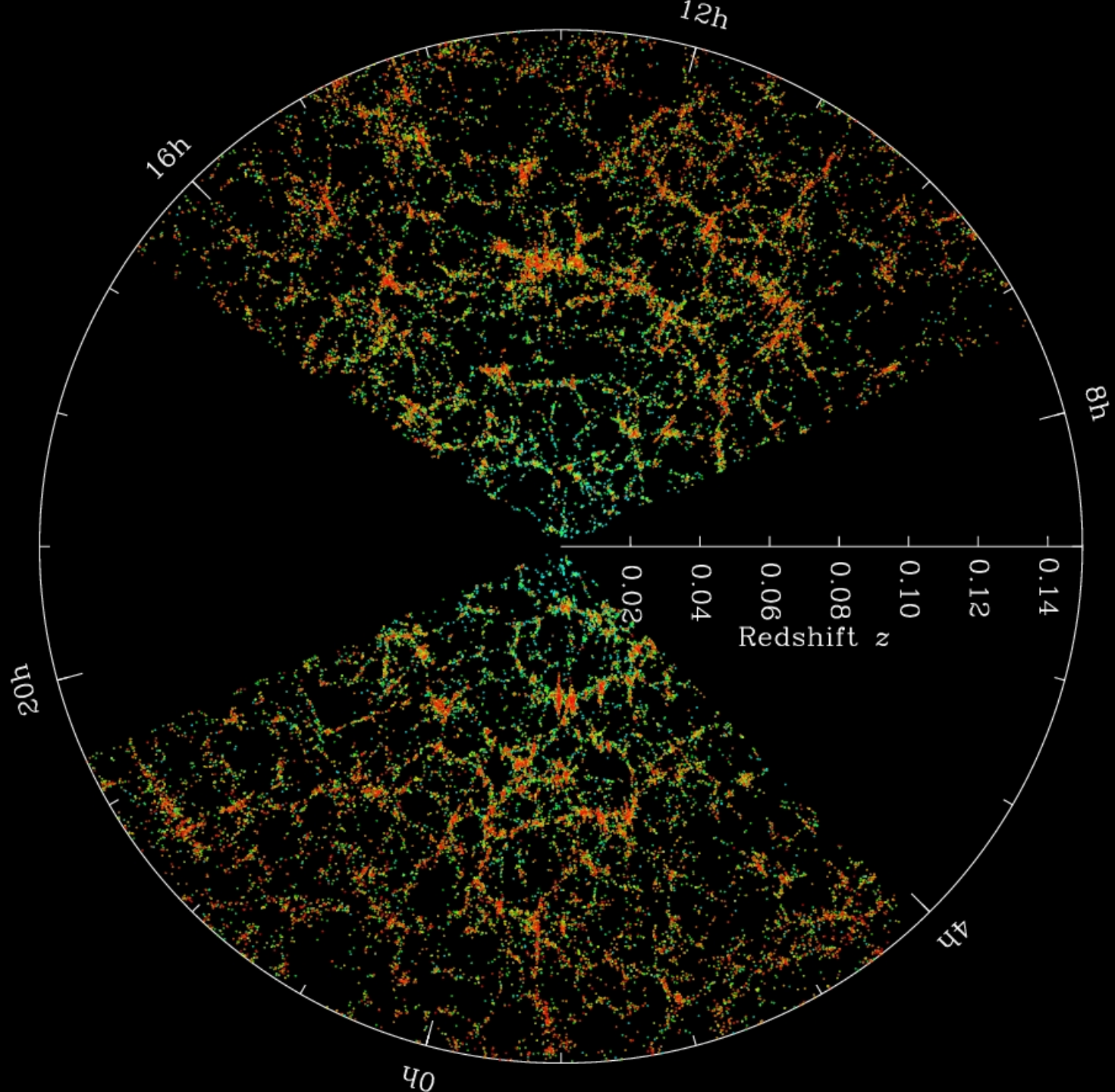
**Gas within/around filament**





125 Mpc/h





Removed (preliminary result)



# Conclusion:

- Most of the baryons in the Universe is in a warm-hot diffuse status, for which X-ray observation is hard to measure.
- We probe gas by cross-correlating the Sunyaev-Zeldovich map from Planck with CFHTLens lensing mass maps and SDSS LRG pair catalogue to probe gas distributions that are difficult to trace.
- Significant correlation is seen with lensing mass. Data is reasonably well fit by a halo model, but requires gas out to  $5 \times$  virial radius. By the virial theorem, the temperature of this gas exactly corresponds to the  $10^5$ — $10^7$ K, i.e. warm-hot intergalactic medium. This is consistent with the finding from numerical simulation.
- We use the aperture photometry filter to the kSZ map, and find the maximum correlation between kSZ-velocity field is at  $\theta=8$  arcmin, corresponding to gas outside virial radius.
- We stack Planck  $y$  on the same sample and find the evidence for a “gas bridge” at the level of  $y = (2.46 \pm 0.93) \times 10^{-8}$  (68% CL). This corresponds to the temperature of gas bridge to be less than  $10^7$  K.
- Our results show that the baryons are no-longer “missing” and they are neither too hot nor too cold, but correlated with underlying mass distribution (lensing and velocity field).

End